

New Characterizations for Bertrand Curves in Minkowski 3-Space

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Abstract: Bertrand curves have been investigated in Lorentzian and Minkowski spaces and some characterizations have been given in [1,2,6]. In this paper, we have investigated the relations between Frenet vector fields and curvatures and torsions of Bertrand curves at the corresponding points in Minkowski 3-space.

Key Words: Bertrand curves, constant curvature and torsion, Minkowski 3- Space.

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§1. Introduction

In the study of the fundamental theory and the characterizations of space curves, the corresponding relations between the curves are the very interesting and important problem. The well-known Bertrand curve is characterized as a kind of such corresponding relation between the two curves. J. Bertrand studied curves in Euclidean 3-space whose principal normals are the principal normals of another curve. Such a curve is nowadays called a Bertrand curve. Bertrand curves have a characteristic property that curvature and torsion are in linear relation. In the recent work [2], the authors studied spacelike and timelike Bertrand curves in Minkowski 3-space. (See also independently obtained results by [6]).

In this paper, we have investigated the relations between Frenet vector fields and curvatures and torsions of Bertrand curves at the corresponding points in Minkowski 3-space.

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§2. Preliminaries

The Minkowski 3-space E_1^3 is the Euclidean 3-space E^3 provided with the standard flat metric given by

$$\langle , \rangle = -dx_1 + dx_2 + dx_3$$

where (x_1, x_2, x_3) is a rectangular coordinate system of E_1^3 . Since \langle , \rangle is an indefinite metric, recall that a vector $v \in E_1^3$ can have one of three Lorentzian causal characters: it can be spacelike if $\langle v, v \rangle > 0$ or $v = 0$, timelike if $\langle v, v \rangle < 0$ and null (lightlike) if $\langle v, v \rangle = 0$ and $v \neq 0$. Similarly, an arbitrary curve $\alpha = \alpha(s)$ in E_1^3 can locally be spacelike, timelike or null (lightlike), if all of its velocity vectors $\alpha'(s)$ are respectively spacelike, timelike or null (lightlike).

Minkowski space is originally from the relativity in Physics. In fact, a timelike curve corresponds to the path of an observer moving at less than the speed of light. Denote by $\{T, N, B\}$ the moving Frenet frame along the curve $\alpha(s)$ in the space E_1^3 . For an arbitrary curve $\alpha(s)$ in the space E_1^3 , the following Frenet formulae are given. If α is timelike curve, then the Frenet formulae read

$$\begin{bmatrix} T' \\ N' \\ B' \end{bmatrix} = \begin{bmatrix} 0 & \kappa & 0 \\ \kappa & 0 & \tau \\ 0 & -\tau & 0 \end{bmatrix} \begin{bmatrix} T \\ N \\ B \end{bmatrix} \quad (1.1)$$

where $\langle T, T \rangle_L = -1, \langle N, N \rangle_L = 1, \langle B, B \rangle_L = 1, \langle T, N \rangle_L = \langle N, B \rangle_L = \langle T, B \rangle_L = 0$. If α is a spacelike curve with a spacelike principal normal, then the Frenet formulae read

$$\begin{bmatrix} T' \\ N' \\ B' \end{bmatrix} = \begin{bmatrix} 0 & \kappa & 0 \\ -\kappa & 0 & \tau \\ 0 & \tau & 0 \end{bmatrix} \begin{bmatrix} T \\ N \\ B \end{bmatrix} \quad (1.2)$$

where $\langle T, T \rangle_L = \langle N, N \rangle_L = 1, \langle B, B \rangle_L = -1, \langle T, N \rangle_L = \langle N, B \rangle_L = \langle T, B \rangle_L = 0$. If α is a spacelike curve with a spacelike binormal, then the Frenet formulae read

$$\begin{bmatrix} T' \\ N' \\ B' \end{bmatrix} = \begin{bmatrix} 0 & \kappa & 0 \\ \kappa & 0 & \tau \\ 0 & \tau & 0 \end{bmatrix} \begin{bmatrix} T \\ N \\ B \end{bmatrix} \quad (1.3)$$

where $\langle T, T \rangle_L = \langle B, B \rangle_L = 1, \langle N, N \rangle_L = -1, \langle T, N \rangle_L = \langle N, B \rangle_L = \langle T, B \rangle_L = 0$ [4,7,11].

§3. Bertrand Curves in Minkowski 3-Space

Definition 3.1([1,2,6]) *Let β_1 and β_2 be two unit speed regular curves in E_1^3 , and $\{T_1, N_1, B_1\}$ and $\{T_2, N_2, B_2\}$ also be Frenet Frames on these curves, respectively. β_1 and β_2 are called Bertrand curves if N_1 and N_2 are linearly dependent. We say that β_2 is a Bertrand mate for β_1 and β_2 are Bertrand curves. And (β_1, β_2) is called a Bertrand couple and we can write*

$$\beta_2(s) = \beta_1(s) + rN_1(s).$$

Theorem 3.1 *If there exists a one-to-one correspondence between the points of the spacelike curves C_1 and C_2 with timelike principal normal, such that at corresponding points P_1 on C_1 and P_2 on C_2 , then the following statements hold:*

- (1) *The curvature κ_1 of C_1 is a constant;*
- (2) *The torsion τ_2 of C_2 is constant;*
- (3) *The unit tangent vector T_1 of C_1 is parallel to the unit tangent vector T_2 of C_2 .*

Then the curve C generated by P that divides the segment P_1P_2 in ratio $m : 1$ is a spacelike Bertrand curve with timelike principal normal.

Proof We shall use the subscripts 1, 2 to designate the geometric quantities corresponding to the curves C_1, C_2 while the same letters without subscripts will refer to the spacelike curve C with timelike principal normal.

Let $\alpha(s), \alpha_1(s), \alpha_2(s)$ be the coordinate vectors at the points P, P_1, P_2 on the curves C, C_1, C_2 respectively. Then the convex combination of points P_1 and P_2 , the equation of point P is

$$\alpha(s) = m\alpha_1(s) + (1 - m)\alpha_2(s), \quad m \in R, \quad (2.1)$$

while by hypothesis,

$$\|T_1\| = \|T_2\| = 1, \quad T_1 = T_2. \quad (2.2)$$

On differentiating Eq.(2.1) we have

$$Tds = mT_1ds_1 + (1 - m)T_2ds_2 = (mds_1 + (1 - m)ds_2)T_1 \quad (2.3)$$

which shows that T is parallel to T_1 and T_2 and always can be chosen so that

$$T = T_1 = T_2, \quad (2.4)$$

and

$$ds = mds_1 + (1 - m)ds_2. \quad (2.5)$$

Differentiating of Eq.(2.4) gives

$$\kappa Nds = \kappa_1 N_1 ds_1 = \kappa_2 N_2 ds_2, \quad (2.6)$$

and if we assume that $\kappa, \kappa_1, \kappa_2$ are positive, then

$$N = N_1 = N_2, \quad (2.7)$$

and

$$\kappa ds = \kappa_1 ds_1 = \kappa_2 ds_2. \quad (2.8)$$

From Eq.(2.4) and Eq.(2.7)

$$B = B_1 = B_2, \quad (2.9)$$

and differentiating

$$\tau Nds = \tau_1 N_1 ds_1 = \tau_2 N_2 ds_2. \quad (2.10)$$

Elimination of ds, ds_1, ds_2 gives

$$\left(\frac{m}{\kappa_1}\right)\kappa + \left(\frac{1-m}{\tau_2}\right)\tau = 1; \kappa_1 \neq 0, \tau_2 \neq 0,$$

which is the desired result, since m, κ_1, τ_2 are constant. If instead of $T_1 = T_2$ were given the condition $B_1 = B_2$, the same result would follow in the same manner. \square

Theorem 3.2 *If condition (c) of Theorem 3.1 is modified so that at corresponding points P_1 and P_2 , the binormals B_1 and B_2 are parallel, then the curve C is a spacelike Bertrand curve with timelike principal normal.*

Proof Since $B_1 = B_2$ then

$$\tau_1 N_1 = \tau_2 N_2 \frac{ds_2}{ds_1}. \tag{2.11}$$

where N_1 and N_2 are the unit normal vectors of α_1 and α_2 at the points P_1 and P_2 with arc-length parametrization. Hence $N_1 = N_2$ and $T_1 = N_1 \times B_1 = N_2 \times B_2$, we know T_1 parallel to T_2 . By Theorem 2.1, C is a spacelike Bertrand curve with timelike principal normal. \square

Theorem 3.3 *If condition (c) of Theorem 3.1 is modified so that at corresponding points P_1 and P_2 the tangent at P_1 is parallel to the binormal B_2 at P_2 , then the curve C is a spacelike Bertrand curve with timelike principal normal.*

Proof Since $T_1 = B_2$, it follow that

$$\kappa_1 N_1 = \tau_2 N_2 \frac{ds_2}{ds_1}. \tag{2.12}$$

Hence N_1 is parallel to N_2 and since N_1 and N_2 are unit vectors,

$$N_1 = N_2 \tag{2.13}$$

and

$$\frac{ds_2}{ds_1} = \frac{\kappa_1}{\tau_2}, \tag{2.14}$$

since $B_1 = T_1 \times N_1 = B_2 \times N_2 = -T_2$, we have

$$B_1 = -T_2. \tag{2.15}$$

Let $\alpha, \alpha_1, \alpha_2$ be the coordinate vectors at the points P, P_1, P_2 on the curves C, C_1, C_2 , respectively. Then

$$\alpha = m\alpha_1 + (1-m)\alpha_2 \tag{2.16}$$

$$\begin{aligned} \frac{d\alpha}{ds} &= m \frac{d\alpha_1}{ds} + (1-m) \frac{d\alpha_2}{ds} \\ \frac{d\alpha}{ds} &= m \frac{d\alpha_1}{ds_1} \frac{ds_1}{ds} + (1-m) \frac{d\alpha_2}{ds_2} \frac{ds_2}{ds} \\ T &= \left(mT_1 + (1-m) \frac{ds_2}{ds_1} T_2 \right) \frac{ds_1}{ds} \\ &= \left(mT_1 + \frac{\kappa_1}{\tau_2} (1-m) T_2 \right) \frac{ds_1}{ds} \end{aligned}$$

$$T = m_1 T_1 + m_2 T_2, \quad (2.17)$$

where

$$\begin{aligned} m_1 &= m \frac{ds_1}{ds} = \frac{m \tau_2}{\sqrt{(m \tau_2)^2 + [\kappa_1(1-m)]^2}}, \quad m_1 = \text{const.} \\ m_2 &= (1-m) \frac{ds_2}{ds} = \frac{(1-m) \kappa_1}{\sqrt{(m \tau_2)^2 + [\kappa_1(1-m)]^2}}, \quad m_2 = \text{const.} \end{aligned}$$

Differentiating Eq.(2.17), one gets

$$\kappa N = \kappa_1 m_1 N_1 \frac{ds_1}{ds} = \kappa_2 m_2 N_2 \frac{ds_2}{ds}. \quad (2.18)$$

Hence

$$N = N_1 = N_2 \quad (2.19)$$

and

$$\kappa = \kappa_1 m_1 \frac{ds_1}{ds} + \kappa_2 m_2 \frac{ds_2}{ds}. \quad (2.20)$$

Using Eq.(2.7) and Eq.(2.9), one finds that

$$B = m_1 B_1 + m_2 B_2. \quad (2.21)$$

Differentiating Eq.(2.11), one gets

$$\tau N = \tau_1 m_1 \frac{ds_1}{ds} N_1 + \tau_2 m_2 \frac{ds_2}{ds} N_2. \quad (2.22)$$

Hence

$$\tau = \tau_1 m_1 \frac{ds_1}{ds} + \tau_2 m_2 \frac{ds_2}{ds}. \quad (2.23)$$

Using Eq.(2.14) and Eq.(2.15), one gets

$$\frac{ds_2}{ds_1} = -\frac{\tau_1}{\kappa_2} = \frac{\kappa_1}{\tau_2}. \quad (2.24)$$

and

$$\frac{\tau_1}{\kappa_1} = -\frac{\kappa_2}{\tau_2}.$$

Let

$$M_1 = m_1 \frac{ds_1}{ds}, \quad M_2 = m_2 \frac{ds_2}{ds}.$$

Then using Eq.(2.20) and Eq.(2.23), one gets

$$\begin{aligned} \frac{\kappa}{M_2 \tau_2} + \frac{\tau}{M_1 \kappa_1} &= \frac{\kappa_1 M_1}{\tau_2 M_2} + \left(\frac{\kappa_2 M_2}{\tau_2 M_2} + \frac{\tau_1 M_1}{\kappa_1 M_1} \right) + \frac{\tau_2 M_2}{\kappa_1 M_1} \\ &= \frac{\kappa_1 M_1}{\tau_2 M_2} + \frac{\tau_2 M_2}{\kappa_1 M_1} = \text{constant}, \quad \frac{\kappa_1}{\tau_2}, \frac{M_1}{M_2} = \text{constant.} \end{aligned}$$

and this is the intrinsic equation of a spacelike Bertrand curve. \square

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